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# Improvement on eight-node quadrilateral element (IQ8) using twice-interpolation strategy for linear elastic fracture mechanics

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ARTICLEINFO	A B S T R A C T
Article history: Received 28 September 2019 Accepted 27 January 2020 Available online 27 January 2020	In this study, an improved eight-node quadrilateral finite element based on a twice-interpolation strategy (TIS) is given for correctly modeling the singular stress field near 2D crack tip of structures. In present approach, the approximation functions for interpolation strategy are established by using the TIS which included nodal values as well as averaged nodal gradients respectively. The stress intensity factors (SIFs) are therefore calculated following the proposed method. The accuracy of the proposed element and its numerical solutions are described by several examples.
Keywords: Twice-interpolation strategy Eight-node quadrilateral element Stress intensity factor	
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## 1. Introduction

Fracture of structures is of great practical as well as theoretical interest. For example, a large number of important engineering structures, such as pressurized aircraft fuselages, ship hulls, storage tanks and pipelines are carefully designed. Concerns about the safety of such structures have led to a large amount of research related to fracture and fatigue of structural materials. The exact numerical model of the cracktip field is still a challenge in the scientific community of computational fracture mechanics. Correct predictions of the singular stress fields near the crack tip are necessary in life prediction, maintenance and safety assessment of advanced structures. Many methods such as analytical, semi-analytical, experimental, numerical methods have been presented to fracture modeling over years. On the basis of numerical methods, it is easy to realize that finite element method (FEM) is a powerful tool because of its wide application to solve many technical problems (Anderssohn, Hofmann, & Bahr, 2018; Bathe, 2006; Bui et al., 2016; Gall, Sehitoglu, & Kadioglu, 1996; Hu et al., 2017; Kuna, 2013; Olgierd Cecil Zienkiewicz, Kenneth Morgan, & Morgan, 2006; Olgierd Cecil Zienkiewicz, Robert L Taylor, Perumal Nithiarasu, & Zhu, 1977; Oliva, Cséplö, Materna, & Bláhová, 1997; Trädegård, Nilsson, & Östlund, 1998). Based on singular crack-tip elements or enriched elements (Denda & Marante, 2004; Duan, Lei, & Li, 2011; Fawkes, Owen, & Luxmoore, 1979; Hu et al., 2017; Jayaswal & Grosse, 1993; Kwon & Akin, 1989; Nash Gifford & Hilton, 1978), etc., the FEM can be used for getting the stress intensity factors. Moreover, a smoothing activity used to stresses recovery at the post-processing state is commonly \* Corresponding author.

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© 2020 Growing Science Ltd. All rights reserved. doi: 10.5267/j.esm.2020.3.005 required by the FEM. Thence, several developments of new or improved numerical techniques were introduced to pass the existing difficulties in the classical methods, for example, the extended finite element method (XFEM) (Bergara, Dorado, Martin-Meizoso, & Martínez-Esnaola, 2017; Feng & Li, 2018; Giner, Sukumar, Tarancón, & Fuenmayor, 2009; Moës, Dolbow, & Belytschko, 1999; Wen & Tian, 2016), meshfree methods and coupled FE-EFG method (Kumar, Singh, & Mishra, 2015; Oliveira et al., 2019; Ramalho, Belinha, & Campilho, 2019; Sun, Hu, & Liew, 2007), smoothed FEM (H. Chen, Wang, Liu, Wang, & Sun, 2016; L. Chen, Liu, Jiang, Zeng, & Zhang, 2011; Liu et al., 2012; Nguyen-Van, Ton-That, Chau-Dinh, & Dao, 2018; Nguyen-Xuan, Liu, Nourbakhshnia, & Chen, 2012; That-Hoang, Nguyen-Van, Chau-Dinh, & Huynh-Van, 2018), extended isogeometric analysis (Bhardwaj, Singh, & Mishra, 2013; Bui, 2015a, 2015b; Shojaee & Daneshmand, 2015; Yin, Yu, Bui, Zheng, & Gu, 2019), phantom-node method (Thanh Chau-Dinh, Mai-Van, Zi, & Rabczuk, 2018; T. Chau-Dinh & Zi, 2011; Thanh Chau-Dinh, Zi, Lee, Rabczuk, & Song, 2012), etc. In this paper, we present an improved eightnode quadrilateral finite element based on a twice-interpolation strategy (TIS) with smooth nodal stresses that can pass the difficulties arose in the standard FEM. In addition to the advantages, FEM also exists some disadvantages such as the discontinuity of gradients of field variables among internal element edges or at nodes. In practice, the post-processing procedure is usually required to get the smoothing operation to the nodal stress. In recent years, some authors (Zheng et al., 2010) presented an improved triangular element for elastostatic problems related to the new concept TIS of the interpolation functions. Their element with various desirable features such as the continuous nodal stress and the higher accuracy of the solutions are not available in the standard elements. The main idea of the TIS is based on the approximation functions that control not only the nodal values but also the averaged nodal gradients as interpolation conditions, see (Zheng et al., 2010) for details. Nevertheless, the major motivation of applying the TIS is to formulate the trial solution as well as its derivatives continuous across inter-element boundaries. The stress by using this strategy can be smoothed over each domain of element to improve the accuracy of solution without the post-processing process. Another important issue should be noted that the TIS does not change the total number of degrees of freedom (DOFs) of the whole system. The TIS is also applied and extended in some references (Bui et al., 2014; Ton-That, Nguyen-Van, & Chau-Dinh, 2019; Ton That, Nguyen-Van, & Chau-Dinh, 2020), respectively.

The main objective of this paper is to improve the eight-node quadrilateral finite element based on TIS for correctly modeling singular stress fields near 2D crack tips of structures. The stress intensity factors (SIFs) calculated by the proposed element are validated against reference solutions. The body of the report is organized into five Sections. In section 2, formulation of this new element (IQ8) for 2D cracks in elastic structures is presented in which the construction of the shape functions and their properties are recommended. The little modification of IQ8 become singular IQ8 around crack tip and the SIFs values are determined in section 3. Several examples are subsequently presented in section 4.

#### 2. Formulation of IQ8 element for linear elastic fracture mechanics

In this section, the construction of the IQ8 shape functions and their properties are briefly given. A point  $\mathbf{x}$  (x, y) is in a eight-node quadrilateral element with eight nodes *i*, *j*, *k*, *m*, *n*, *p*, *q* and *r* as schematically sketched in Fig.1a. We here denote by  $S_{inr}$ ,  $S_{jpn}$ ,  $S_{kqp}$  and  $S_{mrq}$  elements that share nodes *i*, *j*, *k*, *m*, *n*, *p*, *q* and *r*, respectively. All nodes of elements  $S_{inr}$ ,  $S_{jpn}$ ,  $S_{kqp}$  and  $S_{mrq}$  are called the supporting nodes of the point  $\mathbf{x}$  in the IQ8 element. This leads to the support domain for point  $\mathbf{x}$  of IQ8 being larger than the support domain of standard FEM. The trial solution at point  $\mathbf{x}$  is therefore given as

$$u^{h}(\mathbf{x}) = \sum_{l=1}^{n_{sp}} \tilde{N}_{l}(\mathbf{x}) d_{l} = \tilde{\mathbf{N}}(\mathbf{x}) \mathbf{d}$$
(1)

In Eq. (1), the twice-interpolation shape function  $\tilde{N}_l$  is recommended

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$$\tilde{N}_{l} = \underbrace{\phi_{i}N_{l}^{[i]} + \phi_{ix}\overline{N}_{lx}^{[i]} + \phi_{jy}\overline{N}_{ly}^{[i]}}_{\text{node } i} + \underbrace{\phi_{j}N_{l}^{[j]} + \phi_{jx}\overline{N}_{lx}^{[j]} + \phi_{jy}\overline{N}_{ly}^{[j]}}_{\text{node } j} + \underbrace{\phi_{k}N_{lx}^{[k]} + \phi_{kx}\overline{N}_{lx}^{[k]} + \phi_{ky}\overline{N}_{ly}^{[k]}}_{\text{node } k} + \underbrace{\phi_{m}N_{l}^{[m]} + \phi_{mx}\overline{N}_{lx}^{[m]} + \phi_{my}\overline{N}_{ly}^{[m]}}_{\text{node } m} + \underbrace{\phi_{m}N_{lx}^{[n]} + \phi_{mx}\overline{N}_{lx}^{[n]} + \phi_{my}\overline{N}_{ly}^{[n]}}_{\text{node } m} + \underbrace{\phi_{m}N_{lx}^{[n]} + \phi_{mx}\overline{N}_{lx}^{[n]} + \phi_{my}\overline{N}_{ly}^{[n]}}_{\text{node } m} + \underbrace{\phi_{m}N_{lx}^{[m]} + \phi_{mx}\overline{N}_{lx}^{[m]} + \phi_{my}\overline{N}_{ly}^{[m]}}_{\text{node } m} + \underbrace{\phi_{m}N_{lx}^{[m]} + \underbrace{\phi_{m}N_{lx}^{[m]} + \phi_{my}\overline{N}_{ly}^{[m]}}_{\text{node } m} + \underbrace{\phi_{m}N_{lx}^{[m]} + \underbrace{\phi_{m}N_{lx}^{[m]} + \phi_{my}\overline{N}_{ly}^{[m]}}_{\text{nod } m} + \underbrace{\phi_{m}N_{lx}^{[m]} + \underbrace{\phi_{m}N_{lx}^{[m]} + \phi_{my}\overline{N}_{ly}^{[m]}}_{\text{nod } m} + \underbrace{\phi_{m}N_{lx}^{[m]} + \underbrace{\phi_{m$$

in which  $d_f$  and  $N_f^{(i)}$  are called the nodal displacement vector and the shape function related to node *i*. Moreover,  $n_{sp}$  is also called the total number of the supporting nodes due to the point **x**. The formulation of the average derivative of the shape functions at node *i* is presented as below and the same is built for other nodes.

$$\bar{N}_{l,x}^{[i]} = \sum_{e \in S_{inr}} \left( \omega_e N_{l,x}^{[i]/e]} \right); \ \bar{N}_{l,y}^{[i]} = \sum_{e \in S_{inr}} \left( \omega_e N_{l,y}^{[i]/e]} \right)$$
(3)

In Eq. (3), the term  $N_{l,x}^{[i][e]}$  is the derivative of  $N_l^{[i]}$  calculated in element e, and  $\omega_e$  is called the weight function of element  $e \in S_{inr}$ , which is defined by

$$\omega_e = \frac{\Delta_e}{\sum_{\hat{e} \in S_{inr}} \Delta_{\hat{e}}} \quad \text{with } e \in S_{inr}$$
(4)

with  $\Delta_e$  being the area of the element *e*. In Eq. (2), the functions  $\phi_i$ ,  $\phi_{ix}$  and  $\phi_{iy}$  calling the polynomial related to node *i* must be content with the conditions (see Appendix for proving them)

 $\phi_{i}(\mathbf{x}_{l}) = \delta_{il}, \ \phi_{i,x}(\mathbf{x}_{l}) = 0, \ \phi_{i,y}(\mathbf{x}_{l}) = 0, \ \phi_{ix}(\mathbf{x}_{l}) = 0, \ \phi_{ix,x}(\mathbf{x}_{l}) = \delta_{il}, \ \phi_{ix,y}(\mathbf{x}_{l}) = 0$ 

$$\phi_{iy}\left(\mathbf{x}_{l}\right) = 0, \ \phi_{iy,x}\left(\mathbf{x}_{l}\right) = 0, \ \phi_{iy,y}\left(\mathbf{x}_{l}\right) = \delta_{il}$$

in which *l* is from the indices *i*, *j*, *k*, *m*, *n*, *p*, *q* and *r*, and

$$\delta_{il} = \begin{cases} 1 & if \quad i = l \\ 0 & if \quad i \neq l \end{cases}$$
(6)

Note that the above conditions have to be twirled in a similar way to different functions, i.e.,  $\phi_j$ ,  $\phi_{jx}$ ,  $\phi_{jy}$ ,  $\phi_k$ ,  $\phi_{kx}$ ,  $\phi_{ky}$ ,  $\phi_m$ ,  $\phi_{mx}$ ,  $\phi_{my}$ ,  $\phi_n$ ,  $\phi_{nx}$ ,  $\phi_{py}$ ,  $\phi_p$ ,  $\phi_{qy}$ ,  $\phi_q$ ,  $\phi_{qy}$ ,  $\phi_r$ ,  $\phi_{rx}$ , and  $\phi_{ry}$ . These polynomial basis functions  $\phi_i$ ,  $\phi_{xx}$  and  $\phi_{yy}$  for the IQ8 element are given by (7)

$$\begin{split} \phi_{i} &= L_{i} + L_{i}^{2}L_{j} + L_{i}^{2}L_{k} + L_{i}^{2}L_{m} + L_{i}^{2}L_{p} + L_{i}^{2}L_{q} + L_{i}^{2}L_{r} \\ &-L_{i}L_{j}^{2} - L_{i}L_{k}^{2} - L_{i}L_{m}^{2} - L_{i}L_{n}^{2} - L_{i}L_{q}^{2} - L_{i}L_{r}^{2} \\ \phi_{ix} &= -(x_{i} - x_{j})(L_{i}^{2}L_{j} + 0.5L_{i}L_{j}L_{k} + 0.5L_{i}L_{j}L_{m} + 0.5L_{i}L_{j}L_{n} + 0.5L_{i}L_{j}L_{p} + 0.5L_{i}L_{j}L_{q} + 0.5L_{i}L_{j}L_{r}) \\ &-(x_{i} - x_{k})(L_{i}^{2}L_{k} + 0.5L_{i}L_{k}L_{j} + 0.5L_{i}L_{k}L_{m} + 0.5L_{i}L_{k}L_{n} + 0.5L_{i}L_{k}L_{p} + 0.5L_{i}L_{k}L_{q} + 0.5L_{i}L_{k}L_{r}) \\ &-(x_{i} - x_{m})(L_{i}^{2}L_{m} + 0.5L_{i}L_{m}L_{j} + 0.5L_{i}L_{m}L_{k} + 0.5L_{i}L_{m}L_{n} + 0.5L_{i}L_{m}L_{p} + 0.5L_{i}L_{m}L_{q} + 0.5L_{i}L_{m}L_{r}) \\ &-(x_{i} - x_{m})(L_{i}^{2}L_{m} + 0.5L_{i}L_{m}L_{j} + 0.5L_{i}L_{m}L_{k} + 0.5L_{i}L_{m}L_{m} + 0.5L_{i}L_{m}L_{q} + 0.5L_{i}L_{m}L_{q} + 0.5L_{i}L_{m}L_{r}) \\ &-(x_{i} - x_{n})(L_{i}^{2}L_{m} + 0.5L_{i}L_{m}L_{j} + 0.5L_{i}L_{m}L_{k} + 0.5L_{i}L_{m}L_{m} + 0.5L_{i}L_{m}L_{q} + 0.5L_{i}L_{m}L_{q} + 0.5L_{i}L_{m}L_{r}) \\ &-(x_{i} - x_{q})(L_{i}^{2}L_{q} + 0.5L_{i}L_{q}L_{j} + 0.5L_{i}L_{p}L_{m} + 0.5L_{i}L_{p}L_{m} + 0.5L_{i}L_{p}L_{q} + 0.5L_{i}L_{p}L_{q} + 0.5L_{i}L_{p}L_{r}) \\ &-(x_{i} - x_{q})(L_{i}^{2}L_{q} + 0.5L_{i}L_{r}L_{j} + 0.5L_{i}L_{p}L_{m} + 0.5L_{i}L_{p}L_{m} + 0.5L_{i}L_{p}L_{q} + 0.5L_{i}L_{p}L_{q}) \\ &-(x_{i} - x_{j})(L_{i}^{2}L_{q} + 0.5L_{i}L_{r}L_{k} + 0.5L_{i}L_{r}L_{m} + 0.5L_{i}L_{p}L_{m} + 0.5L_{i}L_{p}L_{q} + 0.5L_{i}L_{p}L_{q}) \\ &-(y_{i} - y_{i})(L_{i}^{2}L_{q} + 0.5L_{i}L_{p}L_{k} + 0.5L_{i}L_{p}L_{m} + 0.5L_{i}L_{p}L_{p} + 0.5L_{i}L_{p}L_{q} + 0.5L_{i}L_{p}L_{q}) \\ &-(y_{i} - y_{m})(L_{i}^{2}L_{k} + 0.5L_{i}L_{m}L_{j} + 0.5L_{i}L_{m}L_{k} + 0.5L_{i}L_{m}L_{p} + 0.5L_{i}L_{m}L_{q} + 0.5L_{i}L_{m}L_{q}) \\ &-(y_{i} - y_{m})(L_{i}^{2}L_{m} + 0.5L_{i}L_{m}L_{j} + 0.5L_{i}L_{m}L_{k} + 0.5L_{i}L_{m}L_{m} + 0.5L_{i}L_{m}L_{q} + 0.5L_{i}L_{m}L_{q} + 0.5L_{i}L_{m}L_{q}) \\ &-(y_{i} - y_{m})(L_{i}^{2}L_{m} + 0.5L_{i}L_{m}L_{j} + 0.5L_{i}L_{m}L_{m} + 0.5L_{i}L_{m}L_{m} + 0.5L_{i}L_{m}L$$

(5)



**Fig. 1.** (a) Eight-node quadrilateral element in 2D with twice-interpolation strategy and its support domain, (b) Illustration of the support domains around crack tip.

In Eq. (7), other functions can be also presented in the same way by a circular substitution of indices *i*, *j*, *k*, *m*, *n*, *p*, *q* and *r*. In addition,  $L_i$ ,  $L_j$ ,  $L_k$ ,  $L_m$ ,  $L_p$ ,  $L_q$ , and  $L_r$  are called the area coordinates of the point **x** in the eight-node quadrilateral element *i*, *j*, *k*, *m*, *n*, *p*, *q*, *r* see (Zheng et al., 2010) for more details. It is very important to see these shape functions are sufficient polynomials, hold Kronecker's delta function property and satisfy properties of the partition of unity. The element stiffness matrix **K**<sub>e</sub> can be finally expressed as

$$\mathbf{K}_{e} = \int_{\Omega_{e}} \tilde{\mathbf{B}}_{e}^{T} \mathbf{D} \tilde{\mathbf{B}}_{e} d\Omega$$
(8)

with **D** is elastic tensor and

$$\tilde{\mathbf{B}}_{e} = \begin{bmatrix} \tilde{N}_{1,x} & 0 & \tilde{N}_{2,x} & 0 & \dots & \tilde{N}_{l,x} & 0 & \dots & \tilde{N}_{n_{sp},x} & 0 \\ 0 & \tilde{N}_{1,y} & 0 & \tilde{N}_{2,y} & \dots & 0 & \tilde{N}_{l,y} & \dots & 0 & \tilde{N}_{n_{sp},y} \\ \tilde{N}_{1,y} & \tilde{N}_{1,x} & \tilde{N}_{2,y} & \tilde{N}_{2,x} & \dots & \tilde{N}_{l,y} & \tilde{N}_{l,x} & \dots & \tilde{N}_{n_{sp},y} & \tilde{N}_{n_{sp},x} \end{bmatrix}_{3 \times 2n_{sp}}$$
(9)

 $n_{sp}$  is recalled the total number of the supporting nodes due to the point **x**.

#### 3. Modified IQ8 element around crack tip

For a long time, finite element analyses have been excellent tools for utilization in the evaluation of cracking in structures. In the mid 1970s, (Barsoum, 1974), (Henshell & Shaw, 1975) independently discovered that by taking the mid-side nodes of an element that are adjacent to the tip of crack and moving them to the quarter-point of the element side, the singular stress field which occurs at the tip of crack could be created. It was a meaningful discovery since researchers had spent goodly efforts trying to develop special elements which could capture this behavior. While some of these special elements have achieved some success, implementing these elements into general purpose finite element codes was not practical. The fact that the quarter-point element (QPE) may be used in any finite element code makes it highly valuable.



**Fig. 2.** Eight-node quadrilateral element: (a) with quarter-point nodes, (b) nodal lettering for stress intensity computation.

In this paper, we use quarter-point techniques to modify eight-node quadrilateral element as show in Fig.2a if node 1 is considered as crack tip point. The support domains around crack tip are also presented in Fig.1b, respectively. The SIFs are directly evaluated from

$$K_{I} = \frac{G}{K+1} \sqrt{\left(\frac{2\pi}{L}\right)} \left\{ 4 \left( \dot{v_{B}} - \dot{v_{D}} \right) - \left( \dot{v_{C}} - \dot{v_{E}} \right) \right\}}$$

$$K_{II} = \frac{G}{K+1} \sqrt{\left(\frac{2\pi}{L}\right)} \left\{ 4 \left( \dot{u_{B}} - \dot{u_{D}} \right) - \left( \dot{u_{C}} - \dot{u_{E}} \right) \right\}}$$
(10)

where G is shear modulus, K is 3-4v for plane strain and (3-v)/(1+v) for plane stress, L is QPE length along crack face, and u', v' are local displacement along and normal to crack face as depicted in Fig.2b.

#### 4. Numerical solutions

(a)

In this section, we will examine and assess the IQ8 element through numerical examples.

#### 4.1 An edge crack under tensile load

Firstly, a finite rectangular plate with an edge crack subjected to a uniform tensile load on the top of it is given. The geometry of this test specimen is schematically depicted in Fig.3a. The bottom edge of the specimen is fully clamped. The dimensionless geometric parameters for this specimen is L = 16 for the length of the specimen, W = 7 for the width, and a = W/2 for crack length, respectively.



Fig. 3. Geometric notation of an edge crack under (a) tensile load, (b) shear load and (c) slant edge crack plate under tension.



**Fig. 4.** (a) Convergence of the SIFs for an edge crack under tensile load, (b) Stress field under tensile load, (c) Convergence of the SIFs (KI) for an edge crack under shear load, (d) Convergence of the SIFs (KII) for an edge crack under shear load.

The value  $\sigma = 1$  is assigned to the top of the plate, and only mode I can be considered. The numerical mode-I SIF is validated with the analytical solutions given by (Ewalds & Wanhill, 1989)  $K_I = C\sigma\sqrt{\pi a}$ , where *C* is determined by

$$C = 1.12 - 0.231 \left(\frac{a}{W}\right) + 10.55 \left(\frac{a}{W}\right)^2 - 21.72 \left(\frac{a}{W}\right)^3 + 30.39 \left(\frac{a}{W}\right)^4$$
(11)

The computed results and the convergence of mode-I SIF are thus presented in Fig. 4a, which indicates higher accuracy of the mode-I SIF gained by the IQ8 than the conventional Q8 when increasing the number of elements. To exhibit the performance of IQ8 in stress distribution, we plot the  $\sigma_x$  stress contour which is shown in Fig.4b.

#### 4.2. An edge crack under shear load

By changing from uniform tensile load to uniform shear load on the top of the plate, we have the next example. The geometry of this test specimen is also schematically depicted in Fig.3b. The bottom edge of the specimen is fully clamped. The dimensionless geometric parameter for this specimen is repeated as follows: L = 16 for the length of the specimen, W = 7 for the width, and a = W/2 for crack length. The value  $\tau = 1$  is assigned to the top of the plate. This is a mixed mode problem and the exact solutions are K<sub>I</sub> = 34.0 and K<sub>II</sub> = 4.55 for this case (Moës et al., 1999). Once again, the computer results are presented in Fig. 4c & d, which shows higher accuracy than the conventional Q8 when increasing the number of elements.

#### 4.3 A slant edge crack plate under tension

A rectangular plate is shown in Fig.3c with the dimensions of L = 2.5 cm and b = W = 1.0 cm. An oblique edge crack of length *a* is in plate with  $\beta = 67.5^{\circ}$ . Material properties of the plate are: E = 190GPa, v = 0.25. The stress intensity factors of the crack tip K<sub>I</sub> and K<sub>II</sub> are calculated with the strategic mesh by MATLAB code as shown in Fig.5a while the plate is subjected to uniform tension on the ends. The analytical solution of this model is given as

$$K_{I} = F_{I}\sigma\sqrt{\pi a}$$
  $K_{II} = F_{II}\sigma\sqrt{\pi a}$   $\Rightarrow$   $F_{i} = \frac{K_{i}}{\sigma\sqrt{\pi a}}$  with  $i = I, II$  (11)

Fig.5b show the normalized stress intensity factors (F<sub>I</sub>, F<sub>II</sub>) and the comparison with the results by (Aliabadi et al., 1989) who used the boundary element method to perform similar studies on the same problem. The results are in good agreement.



Fig. 5. (a) The strategic mesh, (b) Normalized SIFs (F<sub>I</sub>, F<sub>II</sub>).

#### 4.4 A three-point bending beam with one edge crack

Finally, let us consider a three-point bending beam with one edge crack and the strategic mesh as depicted in Figs.6 and 7a. The geometric parameters of this structure are W = 6, L = 12 and a = 0.3W. The structure is also subjected to a concentrated force F = 1 and the edge crack is located at the mid-span of its, as the trend of only developing mode I-crack.



Fig. 6. Geometric notation of a three-point bending beam.

The analytical solution with crack length a = 0.3W is given by (Kang et al. 2015; Srawley, 1976). Additionally, higher value of the SIFs computed based on the PUM (Wu & Cai, 2014) over the analytical results can be observed in this reference. The calculated mode-I SIF of this structure based on the IQ8 are compared with the PUM and the analytical method as in Fig.7b with good convergence for the IQ8, respectively.



Fig. 7. (a) The strategic mesh, (b) Mode-I SIF of a three-point bending beam with an edge crack.

### 5. Conclusions

An accurate numerical method based on the IQ8 element is firstly developed for correctly modeling 2D cracks in structures. We have used the IQ8 element to analyze the SIFs of some linear elastic fracture mechanics problems in 2D, respectively. In each case of the study with different load, the SIFs are calculated, and they have been found to agree well with the other analytical results, or with the other numerical methods. Based on the IQ8 element, the present numerical solutions offer higher accuracy than others of the common element. And the applicability of IQ8 element has been clearly shown as above section. The better numerical solutions and the smoother distributions of stresses around the crack tip which are not achieved by the standard elements will be provided when using the IQ8 elements. The computational time of the new IQ8 element was not included in present paper. In fact, this computational time is higher than that based on the standard element because of the TIS procedure, but one will not need a post-processing of any smoothing operation.

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## Appendix

Let us consider a quadrilateral element with eight-node as sketched in Fig.1. The functions  $L_i$ ,  $L_j$ ,  $L_k$ ,  $L_m$ ,  $L_n$ ,  $L_p$ ,  $L_q$  and  $L_r$  are given

$$L_{i} = -\frac{1}{4}(1-\xi)(1-\eta)(1+\xi+\eta); \quad L_{j} = -\frac{1}{4}(1+\xi)(1-\eta)(1-\xi+\eta); \quad L_{k} = -\frac{1}{4}(1+\xi)(1+\eta)(1-\xi-\eta); \quad L_{m} = -\frac{1}{4}(1-\xi)(1+\eta)(1+\xi-\eta)(1+\xi-\eta);$$

$$L_{m} = \frac{1}{2}(1-\xi)(1+\xi)(1-\eta); \quad L_{p} = \frac{1}{2}(1+\xi)(1+\eta)(1-\eta); \quad L_{q} = \frac{1}{2}(1-\xi)(1+\xi)(1+\eta); \quad L_{r} = \frac{1}{2}(1-\xi)(1+\eta)(1-\eta);$$
The device time of the result of the result.

The derivatives of above functions are

$$\begin{cases} \frac{\partial}{\partial \xi} \\ \frac{\partial}{\partial \eta} \end{cases} \begin{bmatrix} L_i & L_j & L_k & L_m & L_n & L_p & L_q & L_r \end{bmatrix} = \\ = \begin{bmatrix} -\frac{1}{4}(-1+\eta)(2\xi+\eta) & \frac{1}{4}(-1+\eta)(\eta-2\xi) & \frac{1}{4}(1+\eta)(2\xi+\eta) & -\frac{1}{4}(1+\eta)(\eta-2\xi) & \xi(-1+\eta) & -\frac{1}{2}(1+\eta)(-1+\eta) & -\xi(1+\eta) & \frac{1}{2}(1+\eta)(-1+\eta) \\ -\frac{1}{4}(-1+\xi)(\xi+2\eta) & \frac{1}{4}(1+\xi)(2\eta-\xi) & \frac{1}{4}(1+\xi)(\xi+2\eta) & -\frac{1}{4}(-1+\xi)(2\eta-\xi) & \frac{1}{2}(1+\xi)(-1+\xi) & -\eta(1+\xi) & -\frac{1}{2}(1+\xi)(-1+\xi) & \eta(-1+\xi) \\ \end{bmatrix}$$
The Level is metric and its inverse are condicitly connected as

The Jacobian matrix and its inverse are explicitly expressed as J =

$$= \begin{bmatrix} -\frac{1}{4}(-1+\eta)(2\xi+\eta) & \frac{1}{4}(-1+\eta)(\eta-2\xi) & \frac{1}{4}(1+\eta)(2\xi+\eta) & -\frac{1}{4}(1+\eta)(\eta-2\xi) & \xi(-1+\eta) & -\frac{1}{2}(1+\eta)(-1+\eta) & -\xi(1+\eta) & \frac{1}{2}(1+\eta)(-1+\eta) \\ -\frac{1}{4}(-1+\xi)(\xi+2\eta) & \frac{1}{4}(1+\xi)(2\eta-\xi) & \frac{1}{4}(1+\xi)(\xi+2\eta) & -\frac{1}{4}(-1+\xi)(2\eta-\xi) & \frac{1}{2}(1+\xi)(-1+\xi) & -\eta(1+\xi) & -\frac{1}{2}(1+\xi)(-1+\xi) & \eta(-1+\xi) \end{bmatrix} \\ \times \begin{cases} x_i & x_j & x_k & x_m & x_n & x_p & x_q & x_r \\ y_i & y_j & y_k & y_m & y_n & y_p & y_q & y_r \end{cases}^{\mathsf{T}}$$

$$\mathbf{J}^{-1} = \frac{1}{\det \mathbf{J}} \begin{bmatrix} J_{3} & J_{4} \end{bmatrix} \quad \text{WIIN}$$

$$J_{1} = -\frac{1}{4} (-1+\xi) (\xi + 2\eta) y_{i} + \frac{1}{4} (1+\xi) (2\eta - \xi) y_{j} + \frac{1}{4} (1+\xi) (\xi + 2\eta) y_{k} - \frac{1}{4} (-1+\xi) (2\eta - \xi) y_{m} + \frac{1}{2} (1+\xi) (-1+\xi) y_{n} - \eta (1+\xi) y_{p} - \frac{1}{2} (1+\xi) (-1+\xi) y_{q} + \eta (-1+\xi) y_{r}$$

$$J_{2} = \frac{1}{4} (-1+\xi) (\xi+2\eta) x_{i} - \frac{1}{4} (1+\xi) (2\eta-\xi) x_{j} - \frac{1}{4} (1+\xi) (\xi+2\eta) x_{k} + \frac{1}{4} (-1+\xi) (2\eta-\xi) x_{m} - \frac{1}{4} (1+\xi) (-1+\xi) x_{n} + \eta (1+\xi) x_{n} + \frac{1}{4} (1+\xi) (-1+\xi) x_{n} - \eta (-1+\xi) x_{n}$$

$$J_{3} = \frac{1}{4}(-1+\eta)(2\xi+\eta)y_{i} - \frac{1}{4}(-1+\eta)(\eta-2\xi)y_{j} - \frac{1}{4}(1+\eta)(2\xi+\eta)y_{k} + \frac{1}{4}(1+\eta)(\eta-2\xi)y_{m} - \xi(-1+\eta)y_{n} + \frac{1}{2}(1+\eta)(-1+\eta)y_{p} + \xi(1+\eta)y_{q} - \frac{1}{2}(1+\eta)(-1+\eta)y_{r}$$

$$J_{4} = -\frac{1}{4}(-1+\eta)(2\xi+\eta)x_{i} + \frac{1}{4}(-1+\eta)(\eta-2\xi)x_{j} + \frac{1}{4}(1+\eta)(2\xi+\eta)x_{k} - \frac{1}{4}(1+\eta)(\eta-2\xi)x_{m} + \frac{1}{4}(1+\eta)(\eta-2\xi)x_{m}$$

$$+\xi(-1+\eta)x_n-\frac{1}{2}(1+\eta)(-1+\eta)x_p-\xi(1+\eta)x_q+\frac{1}{2}(1+\eta)(-1+\eta)x_r$$

## The derivatives of the geometric interpolation functions can be expressed as

$$\begin{aligned} \frac{\partial \phi_i}{\partial L_i} &= 1 + 2L_iL_j + 2L_iL_k + 2L_iL_k + 2L_iL_a + 2L_iL_p + 2L_iL_q + 2L_iL_r - L_j^2 - L_k^2 - L_m^2 - L_n^2 - L_p^2 - L_q^2 - L_r^2 \\ \frac{\partial \phi_i}{\partial L_j} &= L_i^2 - 2L_iL_j \ ; \ \frac{\partial \phi_i}{\partial L_k} = L_i^2 - 2L_iL_k \ ; \ \frac{\partial \phi_i}{\partial L_m} = L_i^2 - 2L_iL_m \ ; \ \frac{\partial \phi_i}{\partial L_n} = L_i^2 - 2L_iL_n \ ; \ \frac{\partial \phi_i}{\partial L_p} = L_i^2 - 2L_iL_p \ ; \ \frac{\partial \phi_i}{\partial L_q} = L_i^2 - 2L_iL_q \ ; \ \frac{\partial \phi_i}{\partial L_r} = L_i^2 - 2L_iL_r \\ \frac{\partial \phi_i}{\partial L_i} &= -(x_i - x_j)(2L_iL_j + 0.5L_jL_k + 0.5L_jL_m + 0.5L_jL_n + 0.5L_jL_q + 0.5L_jL_q + 0.5L_jL_r) - (x_i - x_k)(2L_iL_k + 0.5L_kL_j + 0.5L_kL_m + 0.5L_kL_p + 0.5L_kL_q + 0.5L_kL_r) - (x_i - x_k)(2L_iL_k + 0.5L_kL_p + 0.5L_kL_p + 0.5L_kL_q + 0.5L_qL_q + 0.5L_qL_q - (x_i - x_k)(2L_iL_k + 0.5L_qL_p + 0.5L_qL_q + 0.5L_qL_q + 0.5L_pL_q + 0.5L_qL_q + 0.5L_q$$

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$$\begin{split} \frac{\partial L_{k}}{\partial x} &= (x-x_{k})(1(x-x_{k}))(1(x-x_{k})(x)(1-x_{k})(0)(0)(1,k_{k}) - (x-x_{k})(0)(0)(1,k_{k}) \\ &= (x-x_{k})(0)(0)(1,k_{k}) - (x-x_{k})(0)(0)(1,k_{k}) - ((x-x_{k})(1,k_{k})(1-0)(0)(1,k_{k})) \\ &= \frac{\partial A_{k}}{\partial x_{k}} = -(x-x_{k})(0)(0)(1,k_{k}) - ((x-x_{k})(0)(0)(1,k_{k}) - ((x-x_{k})(0)(0)(1,k_{k})) \\ &= (x-x_{k})(0)(0)(1,k_{k}) - ((x-x_{k})(0)(0)(1,k_{k})) - ((x-x_{k})(0)(0)(1,k_{k})) \\ &= \frac{\partial A_{k}}{\partial x_{k}} = -(x-x_{k})(0)(0)(1,k_{k}) - ((x-x_{k})(0)(0)(1,k_{k})) \\ &= \frac{\partial A_{k}}{\partial x_{k}} = -(x-x_{k})(0)(0)(1,k_{k}) - ((x-x_{k})(0)(0)(1,k_{k})) \\ &= (x-x_{k})(0)(0)(1,k_{k}) - ((x-x_{k})(0)(0)(1,k_{k})) - ((x-x_{k})(0)(0)(1,k_{k})) - ((x-x_{k})(0)(0)(1,k_{k})) \\ &= (x-x_{k})(0)(0)(1,k_{k}) - ((x-x_{k})(0)(0)(1,k_{k})) - ((x-x_{k})(0)(0)(1,k$$

Now, we prove the condition:  $\phi_i(x_i) = \delta_{ii}$ . When i = i, then r = -1; s = -1 and  $L_i = 1$ ,  $L_j = L_k = L_m = L_p = L_p$ =  $L_q = L_r = 0$ , substituting them into above equations we obtain  $\phi_i(x_i) = 1$ . Similarly, when l = j; l = k; l = m; l = n; l = p; l = q or l = r, we respectively obtain  $\phi_i(x_j) = 0$ ;  $\phi_i(x_k) = 0$ ;  $\phi_i(x_m) = 0$ ;

Next, we prove the conditions:  $\phi_{l,v}(x_l) = 0$  and  $\phi_{l,v}(x_l) = 0$ . When  $l \equiv i$ , then r = -1; s = -1 and  $L_i = 1$ ,  $L_j = L_k$  $= L_m = L_p = L_q = L_r = 0$ , substituting them into above equations we have

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$$\begin{cases} \frac{\partial}{\partial x_{1}} \\ \frac{\partial}{\partial x_{1}} \\ \frac{\partial}{\partial x_{2}} \\ \frac{\partial}{\partial x_$$

Similarly, we straightforwardly prove the same for l = j; l = k; l = m; l = n; l = p; l = q or l = r, as well as other conditions.



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